



RUTGERS

Bloch Sphere Representation Of Noisy States

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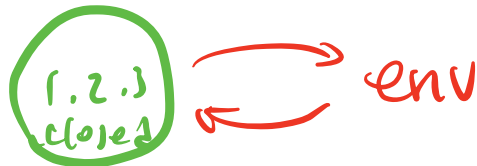
Amey Pimpley

What is Noise?

- Quantum noise is anything that can cause a quantum computer to malfunction.
- In the domain of quantum computing, **noise significantly affects the accuracy of measurements**, which reduces the system's reliability, efficiency, and usefulness.
- When a qubit in a quantum computer gets affected by such kind of noise, the **information in it gets degraded**, just how sound quality is degraded by interference in a phone call. This is called decoherence.

Sources of Quantum Noise

1. **Heisenberg uncertainty:** According to the uncertainty principle, it is impossible to precisely measure both the position and momentum of a particle at the same time. This inherent uncertainty leads to quantum noise in measurements of these quantities.
2. **Noise in quantum states:** The quantum states of particles are inherently probabilistic, which means that they can exhibit random behavior. This randomness can lead to noise in measurements of these states.
3. **Electromagnetic signals from wifi, magnetic field of the earth, etc.**



Decoherence

- In quantum mechanics, **decoherence is the loss of coherence or correlation between the different parts of a quantum system.** In other words, it is the process by which a quantum system **transitions from a state of quantum behavior to classical behavior.**
- This typically happens when a **quantum system interacts with its environment in a way that causes it to become entangled with the surrounding particles,** effectively causing the system to lose its original quantum state.
- Decoherence is an important concept in quantum computing, as it can lead to errors in the calculations performed by a quantum computer.

The Density Matrix

- Matrix representation of possible quantum states weighted by respective probabilities
- Sum of multiplied probabilities and outer products
- Allows us to represent mixed states, such as those produced by noise

$$\rho \equiv \sum_j p_j |\psi_j\rangle \langle \psi_j|$$

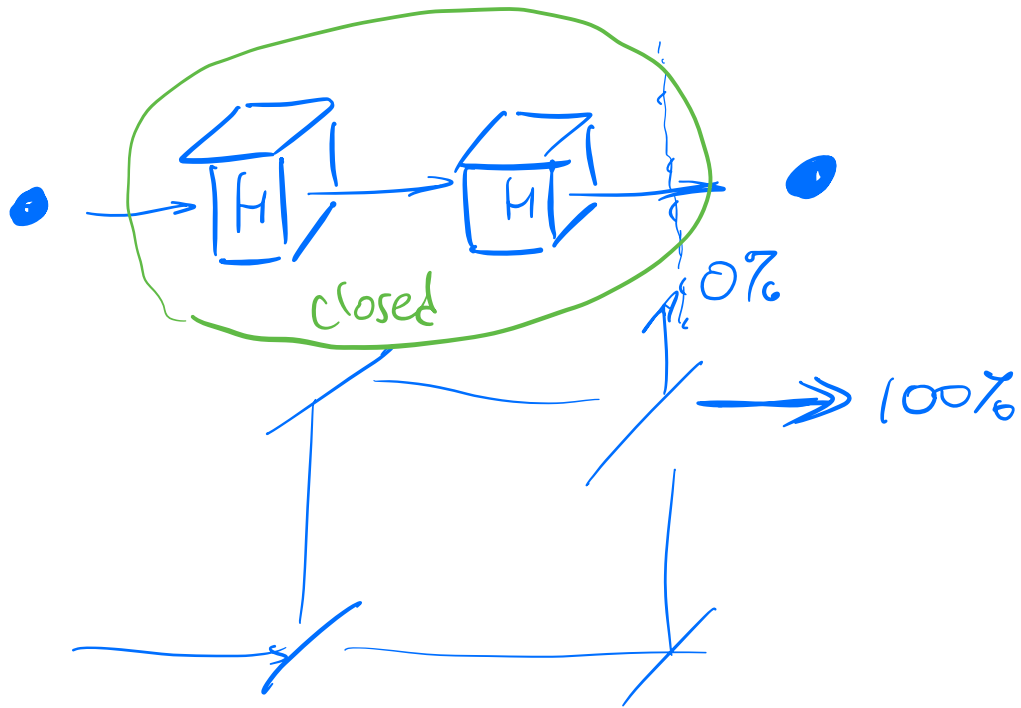
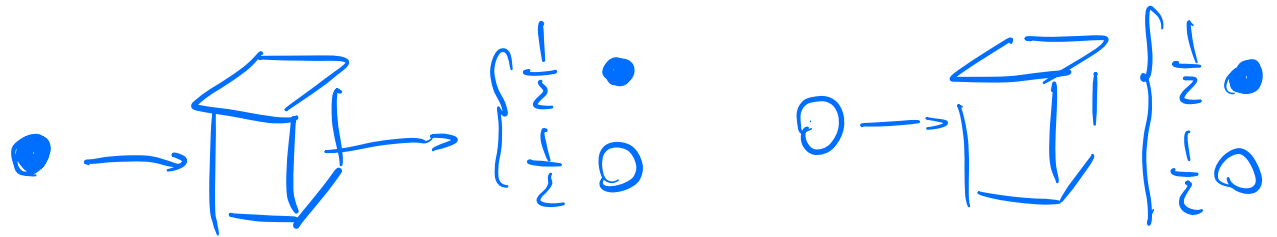
Example: $|0\rangle$ with $\frac{1}{2}$ probability, $|1\rangle$ with $\frac{1}{2}$ probability

$$\begin{aligned} \rho_B &= \frac{1}{2} |0_B\rangle \langle 0_B| + \frac{1}{2} |1_B\rangle \langle 1_B| \\ &= \frac{1}{2} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \end{bmatrix} + \frac{1}{2} \begin{bmatrix} 0 \\ 1 \end{bmatrix} \begin{bmatrix} 0 & 1 \end{bmatrix} \\ &= \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} + \frac{1}{2} \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \\ &= \frac{1}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \end{aligned}$$

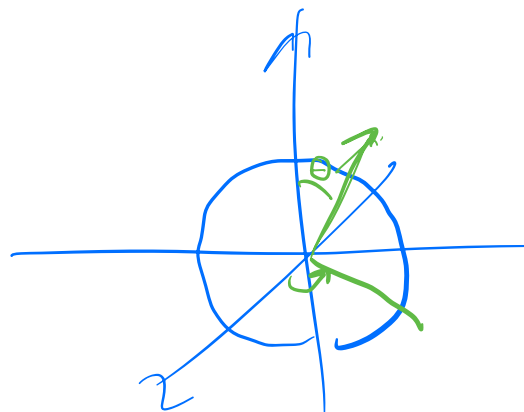
$$\begin{aligned} \rho_{|+\rangle} &= \frac{1}{2} |+\rangle \langle +| \\ &= \frac{1}{2} \begin{bmatrix} \frac{\sqrt{2}}{2} \\ \frac{\sqrt{2}}{2} \end{bmatrix} \begin{bmatrix} \frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} \end{bmatrix} \\ &= \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} \\ &= \frac{I + X}{2} \end{aligned}$$

$$\begin{aligned}
 |+\rangle\langle+| &\xrightarrow{\text{phase flip } \left(\frac{1}{2}\right)} \frac{1}{2}|+\rangle\langle+| + \frac{1}{2}|-\rangle\langle-| \\
 &= \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} + \frac{1}{2} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \\
 &= \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \frac{1}{2} I
 \end{aligned}$$

$$|0\rangle|c_0\rangle \xrightarrow{BF(p:J)} \frac{1}{2}|0\rangle|c_0\rangle + \frac{1}{2}|c_1\rangle$$



$$\alpha|0\rangle + \beta|1\rangle \longrightarrow$$
$$|\alpha|^2 + |\beta|^2 = 1$$



Expansion of the Density Matrix

- In order to map a mixed state on the Bloch sphere using a density matrix, we use this expanded form of ρ , where \vec{r} is the state vector rotated by σ . Solving yields our three coordinates on the sphere.

$$\rho = \frac{1}{2} (\hat{I} + \vec{r} \cdot \hat{\sigma})$$

$$\rho = \frac{1}{2} \hat{I} + \frac{1}{2} r_x \hat{\sigma}_x + \frac{1}{2} r_y \hat{\sigma}_y + \frac{1}{2} r_z \hat{\sigma}_z,$$

where \hat{I} is the identity matrix, and $\hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z$ the X, Y, Z Pauli matrices, respectively:

$$\hat{I} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}, \quad \hat{\sigma}_x = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad \hat{\sigma}_y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}, \quad \hat{\sigma}_z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

Kraus Operators

- The Kraus operators are a **set of mathematical operators that are used to describe the effects of a quantum channel** on a quantum system.
- Each term K_α is a Kraus operator - not necessarily unitary but trace-preserving
 - trace-preserving: $\sum_{\alpha} K_{\alpha} K_{\alpha}^{\dagger} = I$
 - recall that this is similar to unitary operators - unitary channels are the case when there is only one Kraus operator
- Noise channels can be written as a sum of Kraus operators. These can be mathematically derived or produced with functions in Cirq, Qiskit etc.

$$\rho' = \sum_{\alpha} K_{\alpha} \rho K_{\alpha}^{\dagger}$$

Summary of Representations

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

State Vector - simplest, good for pure states

$$\rho \equiv \sum_j p_j |\psi_j\rangle \langle \psi_j|$$

Density Matrix - good for mixed states

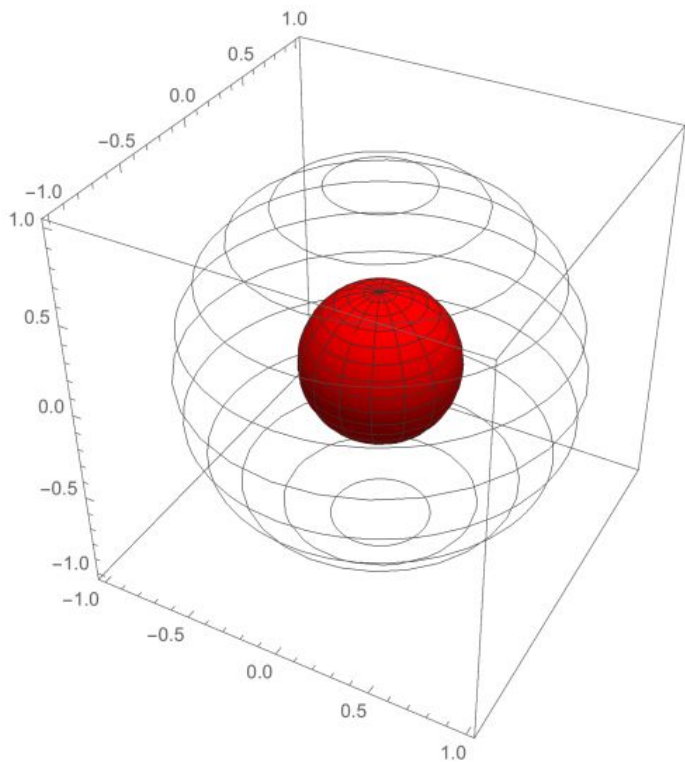
$$\rho = \frac{1}{2} (\hat{I} + \vec{r} \cdot \hat{\sigma})$$

Geometrical expansion of density matrix - good for plotting on sphere

$$\rho' = \sum_{\alpha} K_{\alpha} \rho K_{\alpha}^{\dagger}$$

Kraus operators - describe quantum channels

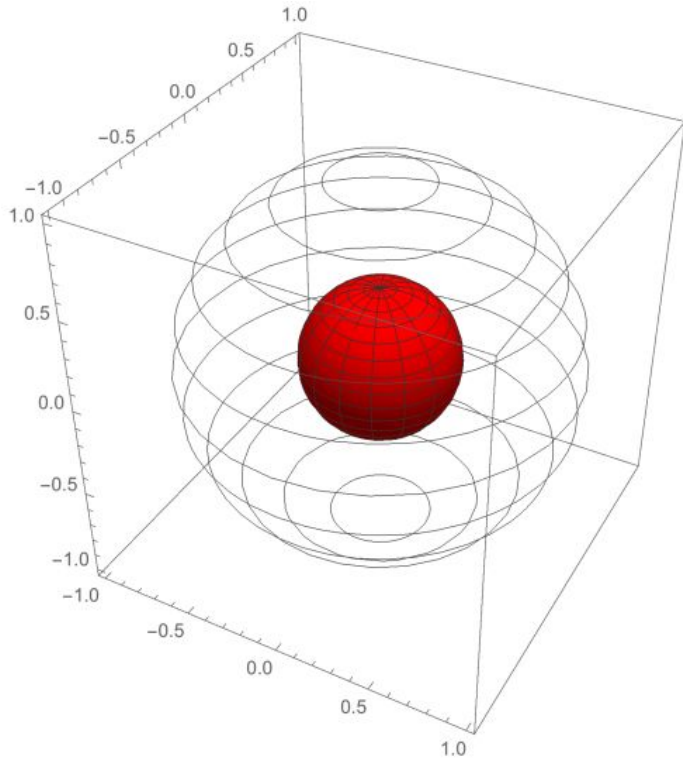
Depolarization



- Takes qubit to maximally mixed state $I/2$ with probability p \rightarrow all info about state is lost
 - maximally mixed state: density matrix is multiple of I
- No effect with probability $1-p$

Effect on Bloch sphere: shrinking toward origin

Depolarization Equations



$$\rho' = p \frac{I}{2} + (1 - p)\rho$$

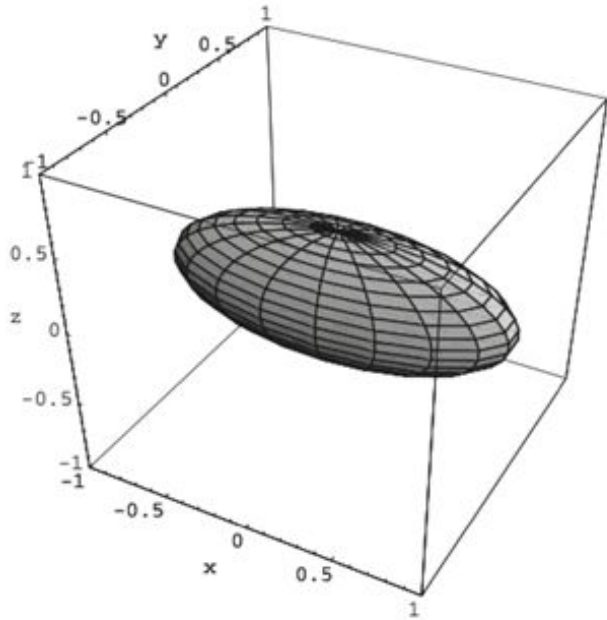
(same as the probability description from last slide)

We can rewrite this to get a Bloch vector:

$$\begin{aligned} \rho' &= p \frac{I}{2} + (1 - p) \frac{1}{2} (I + \vec{r} \cdot \vec{\sigma}) \\ &= p \frac{I}{2} + \frac{(1 - p)}{2} \vec{r} \cdot \vec{\sigma} \\ &= p \frac{I}{2} + \frac{1}{2} (I + \vec{r}' \cdot \vec{\sigma}) \end{aligned}$$

So that $\vec{r}' = (1 - p)\vec{r}$

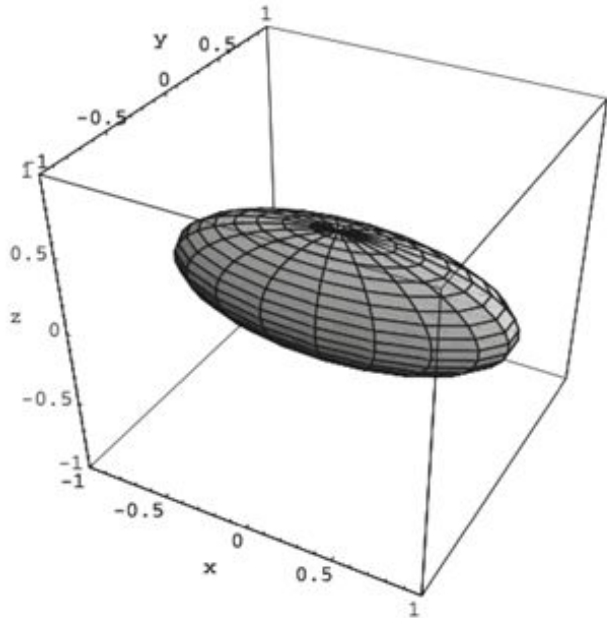
Bit Flip



- Flips bit with probability p
- No effect with probability $1-p$

Effect on Bloch sphere: shrinks across y and z axes

Bit Flip Equations



$$\rho' = p\rho + (1-p)X\rho X$$

$$\rho' = \frac{1}{2}(I + p\vec{r} \cdot \vec{\sigma} + (1-p)X\vec{r} \cdot \vec{\sigma}X)$$

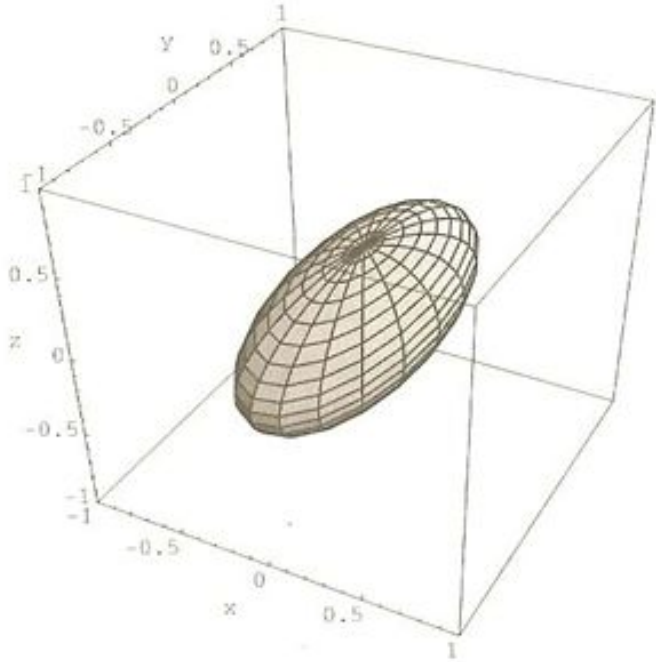
Expanding on the last term gives us:

$$\begin{aligned} X\vec{r} \cdot \vec{\sigma}X &= X(r_xX + r_yY + r_zZ)X \\ &= r_xX - r_yY - r_zZ \end{aligned}$$

Finally giving us:

$$\begin{aligned} \rho' &= \frac{1}{2}(I + r_xX + (2p-1)r_yY + (2p-1)r_zZ) \\ &= \frac{1}{2}(I + \vec{r}' \cdot \vec{\sigma}) \end{aligned}$$

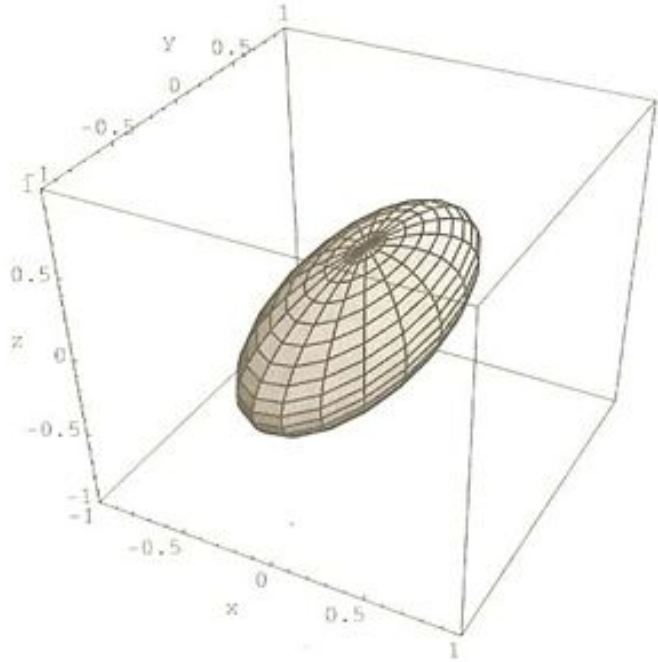
Bit Phase Flip



- Flips bit and sign with probability p
- No effect with probability $1-p$

Effect on Bloch sphere: shrinks across x and z axes

Bit Phase Flip Equations

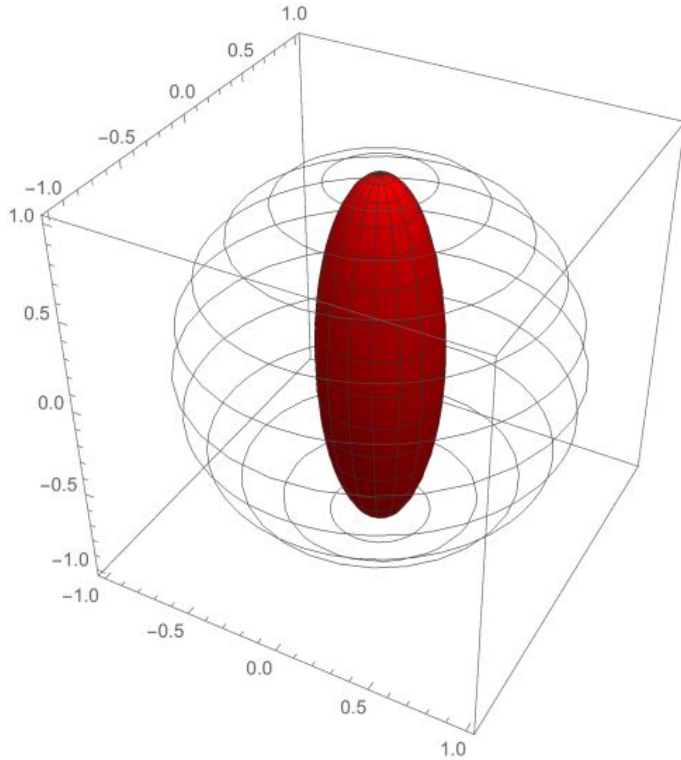


$$\rho' = p\rho + (1 - p)Y\rho Y$$

As we can see, this is identical to bit flip but uses a Y gate. Intuitively we know this must affect only r_x and r_z , but it can be calculated the long way as well.

$$\begin{aligned} \rho' &= \frac{1}{2}(I + (2p - 1)r_x X + r_y Y + (2p - 1)r_z Z) \\ &= \frac{1}{2}(I + \vec{r}' \cdot \vec{\sigma}) \end{aligned}$$

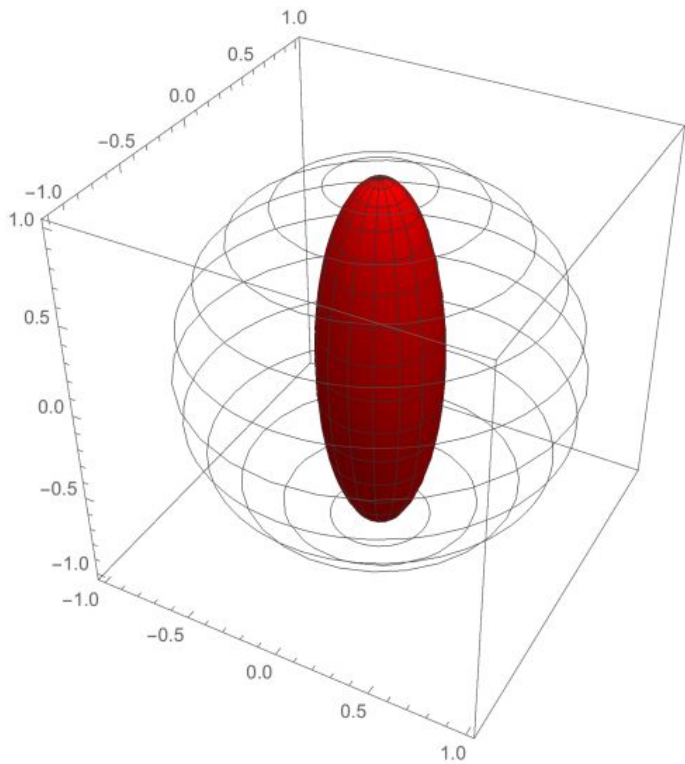
Phase Damping



- Flips sign with probability p
- No effect with probability $1-p$

Effect on Bloch sphere: shrinks across x and y axes

Phase Damping Equations

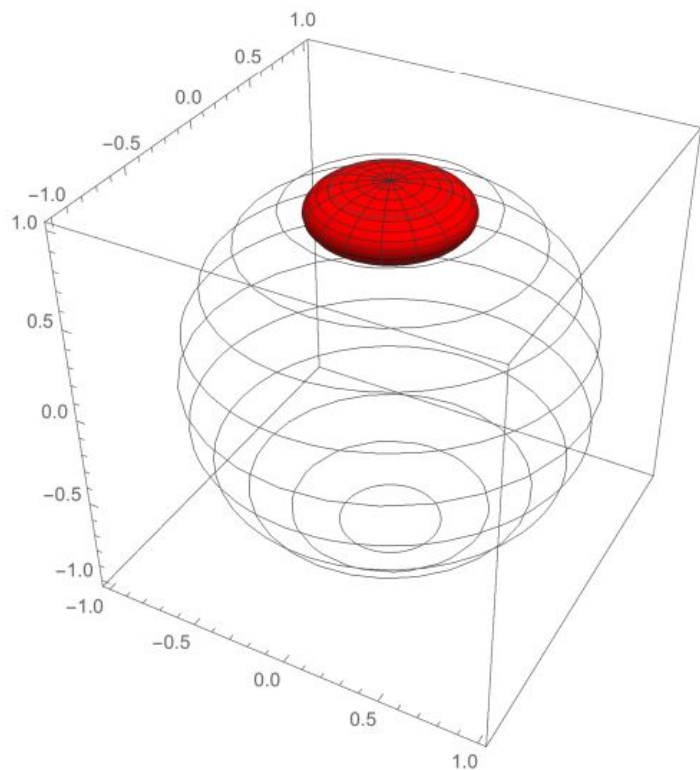


$$\rho' = p\rho + (1 - p)Z\rho Z$$

As we can see, this is identical to bit flip but uses a Z gate. Intuitively we know this must affect only r_x and r_y , but it can be calculated the long way as well.

$$\begin{aligned} \rho' &= \frac{1}{2}(I + (2p - 1)r_x X + (2p - 1)r_y Y + r_z Z) \\ &= \frac{1}{2}(I + \vec{r}' \cdot \vec{\sigma}) \end{aligned}$$

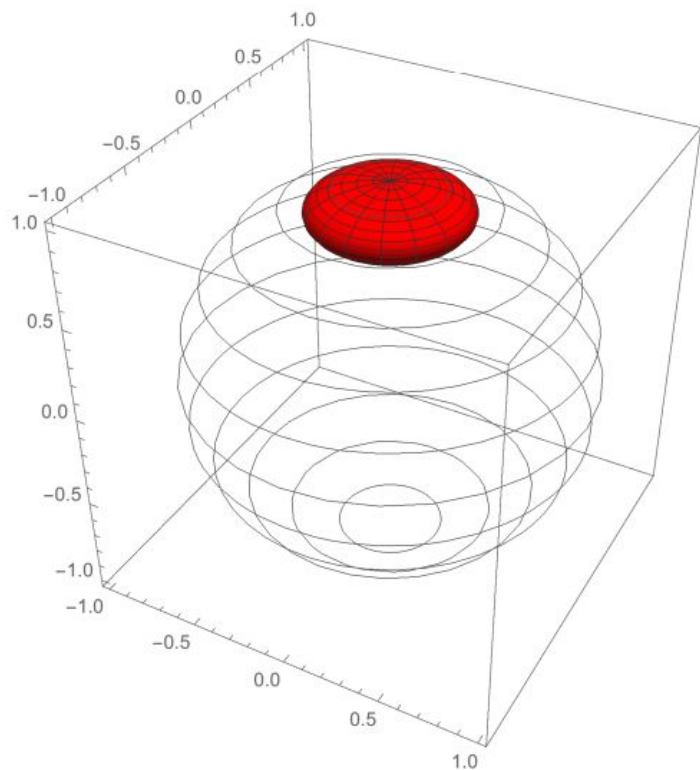
Amplitude Damping



- Models loss of energy
 - $|0\rangle \rightarrow |0\rangle$ with probability 1
 - $|1\rangle \rightarrow |0\rangle$ with probability p
- Can be reversed to model gain of energy

Effect on Bloch sphere: shrinks height and width, center is shifted up or down

Amplitude Damping Equations



Amplitude damping is not unital: it shifts the center of the Bloch sphere, making it more complex

$$\rho' = K_0 \left[\frac{1}{2} (I + \vec{r} \cdot \vec{\sigma}) \right] K_0^\dagger + K_1 \left[\frac{1}{2} (I + \vec{r} \cdot \vec{\sigma}) \right] K_1^\dagger$$

(Expand on board)

$$\rho' = \frac{1}{2} (I + pZ + (\sqrt{1-p})r_x X + (\sqrt{1-p})r_y Y + (1-p)r_z Z)$$

Our Project Goals

- Simulating circuits with different noise channels
 - Visualize final state on Bloch sphere
 - Compare to noiseless circuit results to the ideal one

Questions

References

- Qiskit textbook
<https://qiskit.org/textbook/ch-quantum-hardware/density-matrix.html>
- Cirq documentation https://quantumai.google/cirq/simulate/noisy_simulation
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- Lidar, D. A. (2019). Lecture notes on the theory of open quantum systems. arXiv preprint arXiv:1902.00967